

\vec{E} field of a propagating electromagnetic wave in free-space is given

$$\vec{E} = A \cos(\omega(t - z/c)) \vec{y}$$

Determine \vec{H} for this wave?

(c : speed of wave, t : time, z : position in z -axis, ω : angular frequency)

$$\nabla \times \vec{E} + \frac{d\vec{B}}{dt} = 0 \Rightarrow \frac{d\vec{H}}{dt} = -\frac{1}{\mu_0} \nabla \times \vec{E} \quad (\vec{B} = \mu_0 \vec{H})$$

↳ free-space

where $\nabla \times \vec{E} = \begin{vmatrix} \vec{a}_x & \vec{a}_y & \vec{a}_z \\ d/dx & d/dy & d/dz \\ 0 & E_y & 0 \end{vmatrix} = -\vec{a}_x \frac{dE_y}{dz}$

↳ \vec{E} has only y component

$$\Rightarrow \nabla \times \vec{E} = -\vec{a}_x \frac{d}{dz} (A \cos(\omega(t - z/c))) = -\frac{A\omega}{c} \sin(\omega(t - z/c)) \vec{a}_x$$

$$\Rightarrow \frac{d\vec{H}}{dt} = \frac{A\omega}{\mu_0 c} \sin(\omega(t - z/c)) \vec{a}_x$$

$$\Rightarrow \vec{H} = \frac{A\omega}{\mu_0 c} \int \sin(\omega(t - z/c)) dt \vec{a}_x$$

$$= \frac{A\omega}{\mu_0 c} \frac{-1}{\omega} [\cos(\omega(t - z/c)) + C] \vec{a}_x$$

↳ integral constant which may represent static field if exists which can be ignored

$$\Rightarrow \vec{H} = \frac{-A}{\mu_0 c} \cos(\omega(t - z/c)) \vec{a}_x$$



A simple representation of propagating wave

\vec{E} field of a propagating wave in free-space is given as

$$\vec{E} = \frac{100}{\rho} \sin(\alpha z) \cos(\omega t) \vec{a}_\rho \quad [\text{V/m}] \quad \text{in cylindrical coordinate system.}$$

Determine \vec{H} and α ? (α : constant). No sources exist in the medium ($\rho_f = 0, \vec{J} = 0$)

$$\nabla \times \vec{E} + \frac{d\vec{B}}{dt} = 0, \quad \nabla \times \vec{E} = \frac{1}{\rho} \begin{vmatrix} \vec{a}_\rho & \rho \vec{a}_\phi & \vec{a}_z \\ \frac{d}{d\rho} & \frac{d}{d\phi} & \frac{d}{dz} \\ E_\rho & \rho E_\phi & E_z \end{vmatrix}$$

$$\Rightarrow \nabla \times \vec{E} = \frac{1}{\rho} (-\rho a_\phi \left[-\frac{dE_\rho}{dz} \right]) = +a_\phi \frac{100\alpha}{\rho} \cos(\alpha z) \cos \omega t$$

$$\frac{d\vec{B}}{dt} = \mu_0 \frac{d\vec{H}}{dt} = -\nabla \times \vec{E} = \frac{d\vec{H}}{dt} = \frac{-1}{\mu_0} \frac{100\alpha}{\rho} \cos(\alpha z) \int \cos(\omega t) dt \vec{a}_\phi$$

$$\Rightarrow \vec{H} = \frac{-100\alpha}{\mu_0 \omega \rho} \cos(\alpha z) \sin(\omega t) \vec{a}_\phi$$

Thumb rule is, if \vec{E} & \vec{H} are related to an EM wave, they should satisfy Maxwell's Equations. Then $\nabla \cdot \vec{D} = \rho_f = 0$, $\nabla \cdot \vec{B} = 0$ & $\nabla \times \vec{H} - \frac{d\vec{D}}{dt} = \vec{J} = 0$ (no source)

To determine α , we may check these relations. If one tries $\nabla \cdot \vec{D} = \nabla \cdot \vec{B} = 0$, no knowledge about α can be obtained. Let's utilize Ampere-Maxwell equation.

$$\nabla \times \vec{H} = \frac{d\vec{D}}{dt} = \epsilon_0 \frac{100}{\rho} \sin(\alpha z) \sin(\omega t) (-\omega) \vec{a}_\phi \quad (\equiv \epsilon_0 \frac{d\vec{E}}{dt})$$

$$= \frac{1}{\rho} \begin{vmatrix} \vec{a}_\rho & \rho \vec{a}_\phi & \vec{a}_z \\ \frac{d}{d\rho} & \frac{d}{d\phi} & \frac{d}{dz} \\ 0 & \rho \cdot A_\phi & 0 \end{vmatrix} = \frac{1}{\rho} \left[\vec{a}_\rho \frac{-d(\rho A_\phi)}{dz} + \vec{a}_z \frac{d(\rho A_\phi)}{d\rho} \right]$$

$$\Rightarrow \nabla \times \vec{H} = \frac{-a_\rho}{\rho} \frac{100\alpha^2}{\omega} \sin(\alpha z) \sin(\omega t)$$

$$= -\omega \epsilon_0 \frac{100}{\rho} \sin(\alpha z) \sin(\omega t) \vec{a}_\rho = \epsilon_0 \frac{d\vec{E}}{dt}$$

$$\omega^2 = \omega^2 \epsilon_0 \mu_0 \rightarrow \alpha = \omega \sqrt{\epsilon_0 \mu_0}$$

In vacuum, \vec{B} is observed in cylindrical coordinates as given below;

$$\vec{B} = \begin{cases} B_0 \sin \omega t \vec{a}_z, & \rho < a \\ 0 & \rho > a \end{cases}$$

Determine \vec{E} field? Determine the source?

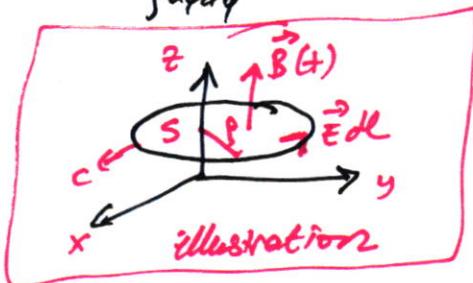
Using Faraday Law, $\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s}$. C : contour attached to open surface S



For $\rho < a$ $\vec{E} \cdot d\vec{\rho} = E_\phi$

$$\int_C \vec{E} \cdot d\vec{l} = E_\phi \cdot \rho \cdot 2\pi = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s}$$

$\rho d\phi \vec{a}_\phi$ $\rho d\rho d\phi \vec{a}_z \rightarrow \vec{B} \parallel d\vec{s}$



$$= -\frac{d}{dt} \int_{\phi=0}^{2\pi} \int_{\rho=0}^{\rho} B_0 \cdot \rho \cdot \sin \omega t \cdot \rho d\rho d\phi$$

$$E_\phi \cdot \rho \cdot 2\pi = \pi \rho^2 B_0 \cdot -d/dt (\sin \omega t) = -\pi \rho^2 B_0 \omega \cos(\omega t)$$

$$\Rightarrow E_\phi = \frac{-\rho}{2} B_0 \omega \cos(\omega t). \text{ In other words } \boxed{\vec{E}_\phi = \frac{-\rho B_0 \omega \cos(\omega t)}{2} \vec{a}_\phi}$$

For $\rho > a$

for $\rho > a, \vec{B} = 0$

$$\oint_C \vec{E} \cdot d\vec{l} = E_\phi \cdot 2\pi \rho = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} = -\pi a^2 B_0 \cos \omega t \cdot \omega$$

$$\Rightarrow \boxed{E_\phi = \frac{-\omega a^2 B_0 \cos \omega t}{2\rho}}$$

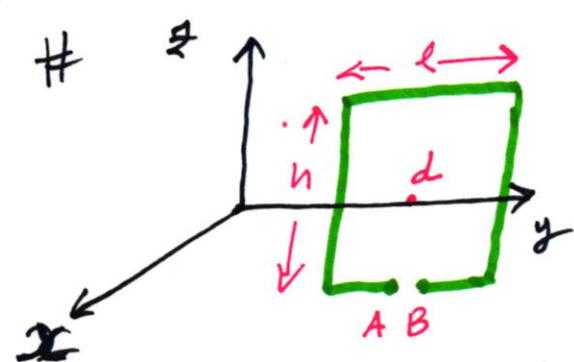
Maxwell eq. to determine sources (ρ_v, \vec{J})

$$\nabla \cdot \vec{D} = \rho_v \rightarrow \nabla \cdot \vec{E} = \rho_v / \epsilon_0 = 0 \Rightarrow \rho_v = 0$$

$$\hookrightarrow \frac{dE_\phi}{d\phi} = 0$$

$$\nabla \times \vec{H} - \frac{d\vec{D}}{dt} = \vec{J}; \quad \nabla \times \vec{H} = \frac{1}{\rho} \begin{vmatrix} \vec{a}_\rho & \rho \vec{a}_\phi & \vec{a}_z \\ \cdot & \cdot & \cdot \\ 0 & 0 & H_z \end{vmatrix} = 0 \Rightarrow \vec{J} = -\frac{d\vec{D}}{dt} = -\epsilon_0 \frac{d\vec{E}}{dt}$$

$$\Rightarrow \vec{J} = \begin{cases} -\epsilon_0 \rho B_0 \omega^2 / 2 \sin(\omega t) \vec{a}_\phi, & \rho < a \\ -\epsilon_0 a^2 \omega^2 B_0 / 2\rho \sin(\omega t) \vec{a}_\phi, & \rho > a \end{cases}$$



An antenna is located in xy plane, symmetrically w.r.t $y=d$ to z & y axis in an electromagnetic field $\vec{B} = B_0 \cos(t - y/c) \vec{a}_x$. Determine the e.m.f between A and B points and the maximum value of it ($l \ll 2c$)

$$E(t) = \text{e.m.f} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s}. \text{ let's choose } \vec{B} \parallel d\vec{s} \text{ so } d\vec{s} = dy dz \vec{a}_x$$

$$\Rightarrow E(t) = -\frac{d}{dt} \int_{y=d-\frac{l}{2}}^{d+\frac{l}{2}} \int_{z=-h/2}^{h/2} B_0 \cos(t - y/c) \vec{a}_x \cdot dy dz \vec{a}_x$$

$$= B_0 \cdot h \cdot \frac{-d}{dt} \int_{d-\frac{l}{2}}^{d+\frac{l}{2}} \cos(t - y/c) dy$$

$$= c B_0 h \frac{d}{dt} \left(\sin(t - y/c) \Big|_{y=d-\frac{l}{2}}^{d+\frac{l}{2}} \right) = ch B_0 \frac{d}{dt} \left[\begin{array}{l} \sin(t - \frac{d+l}{2c}) \\ \sin(t - \frac{d-l}{2c}) \end{array} \right]$$

$$= ch B_0 \left[\cos\left(t - \frac{d+l}{2c}\right) - \cos\left(t - \frac{d-l}{2c}\right) \right]$$

$$= ch B_0 \left[-2 \sin\left(\frac{x+y}{2}\right) \cdot \sin\left(\frac{x-y}{2}\right) \right]$$

$$= -2ch B_0 \sin\left(t - \frac{d}{c}\right) \sin\left(-\frac{l}{2c}\right)$$

$$= 2ch B_0 \sin\left(t - \frac{d}{c}\right) \sin\left(\frac{l}{2c}\right) \quad \frac{l}{2c} \ll 1 \rightarrow \sin\left(\frac{l}{2c}\right) \approx \frac{l}{2c}$$

$$\approx 2ch B_0 \sin\left(t - \frac{d}{c}\right) \frac{l}{2c} \approx B_0 \cdot S \cdot \sin\left(t - \frac{d}{c}\right)$$

↓
area of antenna

Maximum value of $E(t) \rightarrow B_0 \cdot S$ [Volt]

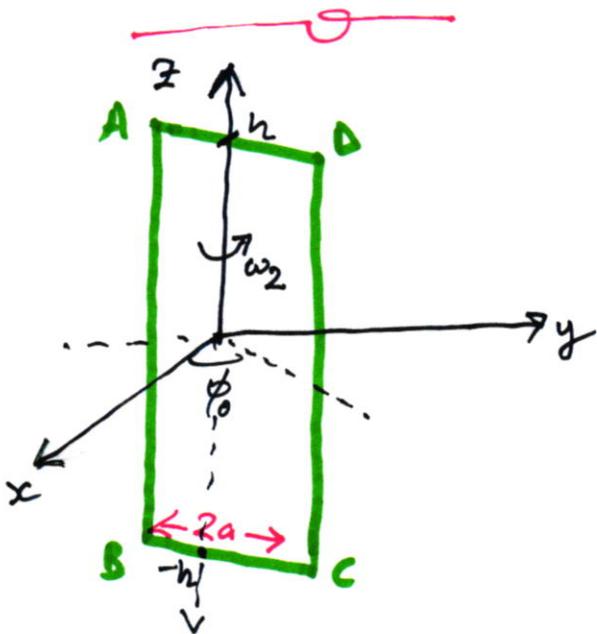
If both \vec{B} and S vary by time; the induced e.m.f becomes

$$\mathcal{E}(t) = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{s} = -\int_S \frac{d\vec{B}}{dt} \cdot d\vec{s} + \oint_C (\vec{v} \times \vec{B}) \cdot d\vec{l}$$

Since $\nabla \times \vec{E} = -\frac{d\vec{B}}{dt}$, the first term of the right side of above equation may be written as $-\int_S \frac{d\vec{B}}{dt} \cdot d\vec{s} = \int_S \nabla \times \vec{E} \cdot d\vec{s} = \oint_C \vec{E} \cdot d\vec{l}$ for which \vec{E} is produced by time varying \vec{B} . Hence;

$$\mathcal{E}(t) = \oint_C [\vec{E} + (\vec{v} \times \vec{B})] \cdot d\vec{l}$$

is the generalized form of Faraday Law (moving surface + time-varying \vec{B})



Consider time-varying $\vec{B} = B_0 \cos \omega_1 t \hat{a}_x$ and a circulating rectangular loop in \hat{a}_ϕ direction with angular speed ω_2 . No charge density ρ exists.

Let's analyze this system. \vec{E} is observed that oscillating in \hat{a}_y direction.

\vec{E} ?, \vec{j} ?, $\mathcal{E}(t)$?

$\phi(t)$?

[At $t=0$, $\phi = \phi_0$]

To determine \vec{E} field, one may utilize $\nabla \cdot \vec{D} = \rho_V$ & $\nabla \times \vec{E} + d\vec{B}/dt = 0$

\vec{E} is observed to be oscillating in \vec{a}_y direction, besides $\rho_V \equiv 0$. Hence

$$\nabla \cdot \vec{E} = \frac{dE_x}{dx} + \frac{dE_y}{dy} + \frac{dE_z}{dz} = 0 \text{ where } E_x = E_z = 0. \text{ So, } \frac{dE_y}{dy} = 0$$

On the other hand; $\nabla \times \vec{E} = d\vec{B}/dt = -\omega_1 B_0 \sin(\omega_1 t) \vec{a}_x$

$$\nabla \times \vec{E} = \begin{vmatrix} \vec{a}_x & \vec{a}_y & \vec{a}_z \\ \frac{d}{dx} & \frac{d}{dy} & \frac{d}{dz} \\ 0 & E_y & 0 \end{vmatrix} = \vec{a}_x (-dE_y/dz) + \vec{a}_z \frac{dE_y}{dx} = +\omega_1 B_0 \sin(\omega_1 t) \vec{a}_x$$

$$\Rightarrow -dE_y/dz \cdot \vec{a}_x = +\omega_1 B_0 \sin(\omega_1 t) \vec{a}_x$$

$$\Rightarrow \boxed{E_y = -\omega_1 B_0 z \sin(\omega_1 t) + C(t)}$$

integral constant
varying by time with
no position dependence
It may be ignored in
calculations -

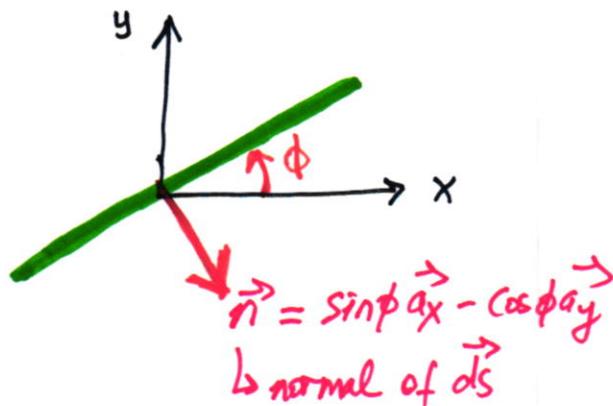
To determine \vec{J} , one may utilize $\nabla \times \vec{H} - \frac{d\vec{D}}{dt} = \vec{J}$

$\nabla \times \vec{H} = 0$ (\vec{H} or \vec{B} is independent of position)

$$-\frac{d\vec{D}}{dt} = -\epsilon_0 \frac{d\vec{E}}{dt} = \boxed{\epsilon_0 \omega_1^2 B_0 z \cos(\omega_1 t) \vec{a}_y = \vec{J}}$$

Let's calculate the magnetic flux through the loop as a function of time, that is $\phi(t)$

$$\phi(t) = \int \vec{B} \cdot d\vec{s}$$



$$\int \vec{B} \cdot d\vec{s} = \int B_0 \cos\omega_1 t \vec{a}_x \cdot (\sin\phi \vec{a}_x - \cos\phi \vec{a}_y) ds$$

$$= B_0 \cos\omega_1 t \sin\phi \cdot A$$

↳ surface area of loop

$$\phi = \phi_0 + \omega_2 t$$

$$\phi(t) = B_0 \cos(\omega_1 t) \sin(\phi_0 + \omega_2 t) \cdot A$$

$$[T \cdot m^2 \cdot Wb]$$

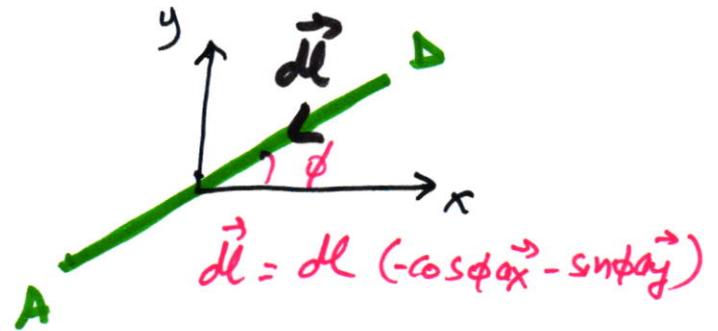
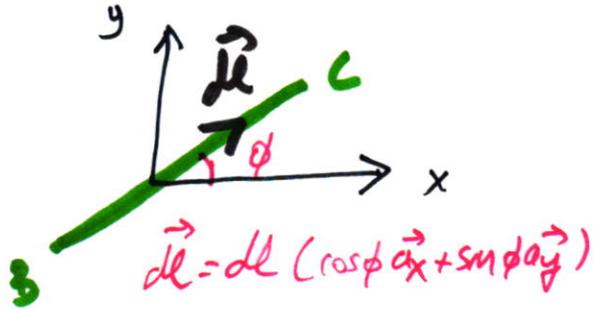
Let's calculate $\oint \vec{E} \cdot d\vec{l}$ on the loop which is the circulation of \vec{E} through the surface.

Note that \vec{E} has only y component for which the $\oint \vec{E} \cdot d\vec{l}$ becomes 0 on AB & CD paths (since $d\vec{l} = dl \hat{a}_x$)

Hence

$$\oint \vec{E} \cdot d\vec{l} = \int_{AB} \vec{E} \cdot d\vec{l} + \int_{CD} \vec{E} \cdot d\vec{l} + \int_{BC} \vec{E} \cdot d\vec{l} + \int_{DA} \vec{E} \cdot d\vec{l} = \int_{BC} \vec{E} \cdot d\vec{l} + \int_{DA} \vec{E} \cdot d\vec{l}$$

for the chosen direction of dS ;



So;

$$\oint \vec{E} \cdot d\vec{l} = \int_{y=-a}^a E \hat{a}_y \cdot dy (\cos \phi \hat{a}_x + \sin \phi \hat{a}_y) + \int_{y=a}^{-a} E \hat{a}_y \cdot (-\cos \phi \hat{a}_x - \sin \phi \hat{a}_y) dy$$

$$= E \sin \phi 2a \Big|_{z=-h}^{z=h} \Big|_{\phi=\phi_0+\omega_2 t} + -E \sin \phi 2a \Big|_{z=h}^{z=-h} \Big|_{\phi=\phi_0+\omega_2 t}$$

$$= B_0 h \omega_1 2a \sin(\phi_0 + \omega_2 t) \sin(\omega_1 t) + B_0 h \omega_1 \sin(\omega_1 t) 2a \sin(\phi_0 + \omega_2 t)$$

$$= 4ah B_0 \omega_1 \sin(\omega_1 t) \sin(\phi_0 + \omega_2 t)$$

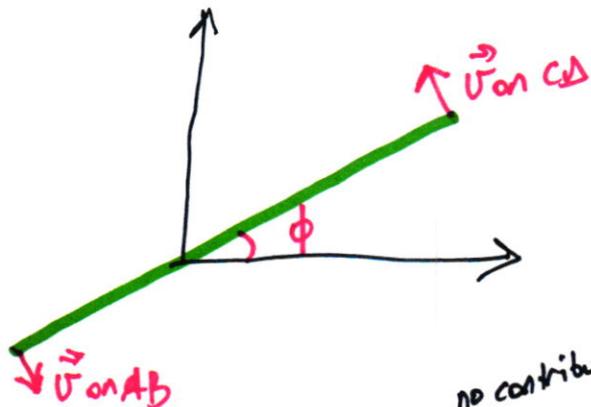
$$\oint \vec{E} \cdot d\vec{l} = \underbrace{4ah}_{A} B_0 \omega_1 \sin(\omega_1 t) \sin(\phi_0 + \omega_2 t) \quad [\text{Volt}]$$

Let's calculate $\oint_C (\vec{v} \times \vec{B}) \cdot d\vec{\ell}$. Note that $\vec{v} \times \vec{B} \parallel \hat{a}_z$ so the contour integral
 $(\hat{a}_x \times \hat{a}_y)$

C is composed of AB & CD paths (z direction)

on AB $d\vec{\ell} = dl (-\hat{a}_z)$ and on CD, $d\vec{\ell} = dl \hat{a}_z$ for chosen dS

on AB, $\vec{v} = \omega_2 \cdot a \cdot \hat{a}_\phi$ and on CD, $\vec{v} = \omega_2 \cdot a \cdot (-\sin\phi \hat{a}_x + \cos\phi \hat{a}_y)$
 $= \omega_2 \cdot a \cdot (\sin\phi \hat{a}_x - \cos\phi \hat{a}_y)$



no contribution

$$\text{So, } \oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = \int_{AB} [\omega_2 a (\sin\phi \hat{a}_x - \cos\phi \hat{a}_y) \times B_0 \cos(\omega_1 t) \hat{a}_x] \cdot d\vec{\ell} (-\hat{a}_z)$$

no contribution +

$$+ \int_{CD} [\omega_2 a (-\sin\phi \hat{a}_x + \cos\phi \hat{a}_y) \times B_0 \cos(\omega_1 t) \hat{a}_x] \cdot d\vec{\ell} (\hat{a}_z)$$

Path integral dl contributes multiplier $2h \int_{z=-h}^h dl = 2h$

$$\text{Hence } \oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = -2h \omega_2 a B_0 \cos(\omega_1 t) \cos\phi \hat{a}_z + -2h \omega_2 a B_0 \cos(\omega_1 t) \cos\phi \hat{a}_z$$

$$= -4ah \omega_2 B_0 \cos(\omega_1 t) \cos\phi \hat{a}_z$$

$$\oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = -A \omega_2 B_0 \cos(\omega_1 t) \cos(\phi_0 + \omega_2 t) \hat{a}_z \text{ Volt}$$

In total

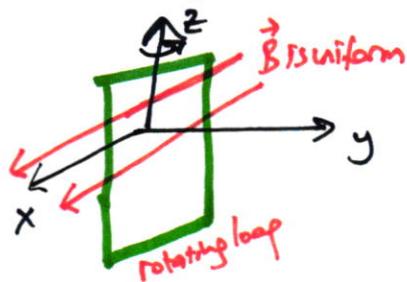
$$\oint_C (\vec{E} + \vec{v} \times \vec{B}) \cdot d\vec{\ell} = A B_0 \omega_1 \sin(\omega_1 t) \sin(\phi_0 + \omega_2 t) - A B_0 \omega_2 \cos(\omega_1 t) \cos(\phi_0 + \omega_2 t)$$

$$\stackrel{?}{=} -d\phi/dt$$

$$\phi = A B_0 \cos(\omega_1 t) \sin(\phi_0 + \omega_2 t) \Rightarrow \frac{d\phi}{dt} = \oint_C (\vec{E} + \vec{v} \times \vec{B}) \cdot d\vec{\ell}$$

If \vec{B} is constant in time, that is $\omega_1 = 0$, $\vec{B} = B_0 \cdot a_x \hat{x}$

Flux through the surface becomes; $\phi = A B_0 \sin(\phi_0 + \omega_2 t)$



Since $\frac{dB}{dt} = 0$, no \vec{E} field is produced

$$\text{so } \int \vec{E} \cdot d\vec{\ell} = 0$$

$$\mathcal{E}(t) = -\frac{d\phi}{dt} = -A B_0 \omega_2 \cos(\phi_0 + \omega_2 t)$$

$$= \oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = -A \omega_2 B_0 \cos(\phi_0 + \omega_2 t)$$

If loop does not rotate, that is $\omega_2 = 0$ while \vec{B} varies by time

Flux through the surface becomes, $\phi = A B_0 \cos(\omega_1 t) \sin \phi_0$

Since, no moving loop $\oint (\vec{v} \times \vec{B}) \cdot d\vec{\ell} = 0$ but $\int \vec{E} \cdot d\vec{\ell} = A B_0 \omega_1 \sin(\omega_1 t) \sin(\phi_0)$

$$= -\frac{d\phi}{dt} = \mathcal{E}(t)$$

If both \vec{B} and loop is constant in time, that is $\omega_1 = \omega_2 = 0$

$\int \vec{E} \cdot d\vec{\ell} = \int \vec{v} \times \vec{B} \cdot d\vec{\ell} = 0$ and no e.m.f is generated

$\phi = \int \vec{B} \cdot d\vec{s} = A B_0 \sin \phi_0 \rightarrow -d\phi/dt = 0$ and no e.m.f is produced

~ THE END ~

Good Luck with your final exams

Yüzyen

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in the preparation of lecture notes"