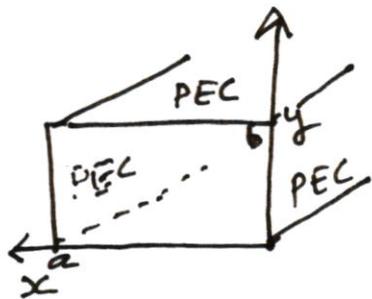


## Rectangular Waveguides (TE, TM Modes are possible)



### TE Modes ( $E_z = 0, H_z \neq 0$ )

$$(\nabla_t^2 + k_c^2) H_z = 0 = (\nabla_t^2 + k_c^2) h_z(x, y) e^{-i\beta z} = 0$$

Now let's separate  $h_z(x, y)$  in terms of multiplication as;

$$h_z(x, y) = X(x) Y(y)$$

and solve differential equation  $(\nabla_t^2 + k_c^2) X(x) Y(y) = 0$

$$Y \frac{d^2 X}{dx^2} + X \frac{d^2 Y}{dy^2} + k_c^2 X Y = 0 = \underbrace{\frac{1}{X} \frac{d^2 X}{dx^2}}_{f(x)} + \underbrace{\frac{1}{Y} \frac{d^2 Y}{dy^2}}_{g(y)} + k_c^2$$

$f(x)$  and  $g(y)$  are functions of  $x$  and  $y$ , respectively. They should be constant since  $k_c$  is constant. Hence, one may write;

$$\left. \begin{array}{l} f(x) + g(y) = -k_c^2 \\ \downarrow \quad \downarrow \\ -k_x^2 - k_y^2 = -k_c^2 \end{array} \right\} k_x^2 + k_y^2 = k_c^2$$

$$f(x) = -k_x^2 \Rightarrow \frac{1}{X} \frac{d^2 X}{dx^2} = -k_x^2 \Rightarrow \frac{d^2 X}{dx^2} + X k_x^2 = 0$$

$$g(y) = -k_y^2 = \frac{1}{Y} \frac{d^2 Y}{dy^2} = -k_y^2 \Rightarrow \frac{d^2 Y}{dy^2} + Y k_y^2 = 0$$

Then,

$$X = A \cos(k_x x) + B \sin(k_x x)$$

$$Y = C \cos(k_y y) + D \sin(k_y y)$$

$$h_z(x, y) = X(x)Y(y) = [A \cos(k_x x) + B \sin(k_x x)] [C \cos(k_y y) + D \sin(k_y y)]$$

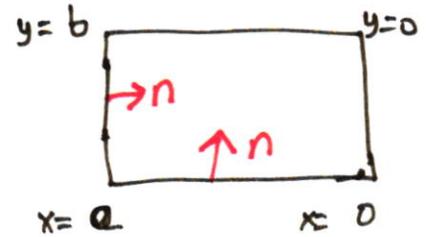
$$H_z(x, y, z) = h_z(x, y) e^{-i\beta z}$$

$A, B, C, D$  constants can be determined ( $k_x$  and  $k_y$  as well) by Boundary conditions. Recall that  $\frac{d\psi}{dn} = 0$  on PEC where  $\psi = H_z$ .

on PEC

$$\frac{d\psi}{dn} = 0 \rightarrow \left. \frac{dh_z}{dx} \right|_{x=0} = 0 \quad (1) \quad \left. \frac{dh_z}{dx} \right|_{x=a} = 0 \quad (2)$$

$$\left. \frac{dh_z}{dy} \right|_{y=0} = 0 \quad (3) \quad \left. \frac{dh_z}{dy} \right|_{y=b} = 0 \quad (4)$$



$$\frac{dh_z}{dx} = Y [-A k_x \sin(k_x x) + B k_x \cos(k_x x)] = 0$$

From (1),  $B=0$ ; from (2)  $k_x = \frac{m\pi}{a}$

$$\frac{dh_z}{dy} = X [-C k_y \sin(k_y y) + D k_y \cos(k_y y)] = 0$$

From (3),  $D=0$ ; from (4)  $k_y = \frac{n\pi}{b}$

So; <sup>arbitrary constant</sup>

$$h_z(x, y) = A.C. \cos\left(\frac{m\pi}{a}x\right) \cos\left(\frac{n\pi}{b}y\right)$$

$m = 0, 1, 2, \dots$  or  $m = 1, 2, \dots$ ,  $m = n = 0$  is not allowed  
 $n = 1, 2, \dots$  or  $n = 0, 1, \dots$

$$\beta = \sqrt{k^2 - k_c^2} \text{ where } k_c^2 = k_x^2 + k_y^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2$$

$$\beta_{mn} = \sqrt{k^2 - \left[ \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 \right]} \rightarrow \text{propagation constant of TE}_{mn} \text{ mode}$$

Each mode has its own cut-off frequency. ( $\beta_{mn} = 0$ )

$$\beta_{mn} = 0 \rightarrow k^2 = k_c^2 = \left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2 \text{ where } k = \frac{2\pi f}{v_p}$$

hence  $f_{c_{mn}} = \frac{c}{2\pi\sqrt{\epsilon_r\mu_r}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$  → cut-off frequency of the  $TE_{mn}$  mode.

At a given frequency  $f$  →  $f > f_c \rightarrow$  propagating mode  
 →  $f = f_c \rightarrow$  cut-off condition  
 →  $f < f_c \rightarrow$  evanescent mode

Generally  $a > b$  is chosen for waveguide applications  Then,

$$f_{c_{10}}^{TE} = \frac{c}{2a\sqrt{\epsilon_r\mu_r}}$$

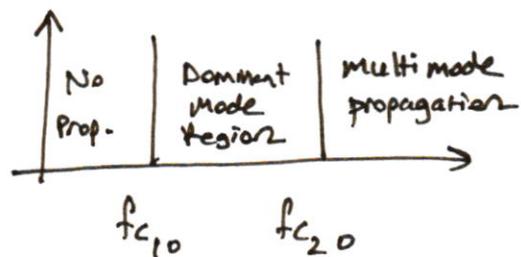
$$f_{c_{20}}^{TE} = \frac{c}{a\sqrt{\epsilon_r\mu_r}}$$

$$f_{c_{01}}^{TE} = \frac{c}{2b\sqrt{\epsilon_r\mu_r}}$$

$$f_{c_{11}}^{TE} = \frac{c}{2} \sqrt{\frac{1}{a^2} + \frac{1}{b^2}} \frac{1}{\sqrt{\epsilon_r\mu_r}}$$

If  $a > 2b$

$$f_{c_{10}} < f_{c_{20}} < f_{c_{01}} < f_{c_{11}}$$



The mode with the lowest cut-off frequency is the "dominant mode".

For a given geometry  $a > b$ ,  $TE_{10}$  is the dominant mode.

## TM Modes ( $E_z \neq 0, H_z = 0$ )

$$(\nabla_t^2 + k_c^2) E_z = 0 = (\nabla_t^2 + k_c^2) e_z(x, y) e^{-i\beta z}$$

Letting  $e_z(x, y) = X(x) Y(y)$  we may write

$$Y \frac{d^2 X}{dx^2} + X \frac{d^2 Y}{dy^2} + X Y k_c^2 = 0$$

$$\Rightarrow \underbrace{\frac{1}{X} \frac{d^2 X}{dx^2}}_{-k_x^2} + \underbrace{\frac{1}{Y} \frac{d^2 Y}{dy^2}}_{-k_y^2} + k_c^2 = 0 \rightarrow k_c^2 = k_x^2 + k_y^2$$

which gives;

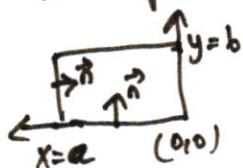
$$\frac{1}{X} \frac{d^2 X}{dx^2} + k_x^2 = 0 \quad ; \quad \frac{1}{Y} \frac{d^2 Y}{dy^2} + k_y^2 = 0$$

Solution of differential equations are;

$$X = A \cos(k_x x) + B \sin(k_x x)$$

$$Y = C \cos(k_y y) + D \sin(k_y y)$$

Boundary condition is that  $\varphi = 0$  on PEC where  $\varphi = E_z$



$$\text{At } x=0, e_z=0 \rightarrow \boxed{A=0}$$

$$\text{At } x=a, e_z=0 \rightarrow \boxed{k_x = \frac{m\pi}{a}}$$

$$\text{At } y=0, e_z=0 \rightarrow \boxed{C=0}$$

$$\text{At } y=b, e_z=0 \rightarrow \boxed{k_y = \frac{n\pi}{b}}$$

So, for TM modes;

$$E_z(x, y, z) = B.D. \sin\left(\frac{m\pi}{a} x\right) \sin\left(\frac{n\pi}{b} y\right) e^{-i\beta_{mn} z}$$

where

$$\beta_{mn} = \sqrt{k^2 + k_c^2} = \sqrt{k^2 - k_x^2 - k_y^2} = \sqrt{k^2 - \left(\frac{m\pi}{a}\right)^2 - \left(\frac{n\pi}{b}\right)^2}$$

For TM modes  $m$  or  $n$  cannot be 0.  $TM_{10}$  or  $TM_{01}$  does not exist.

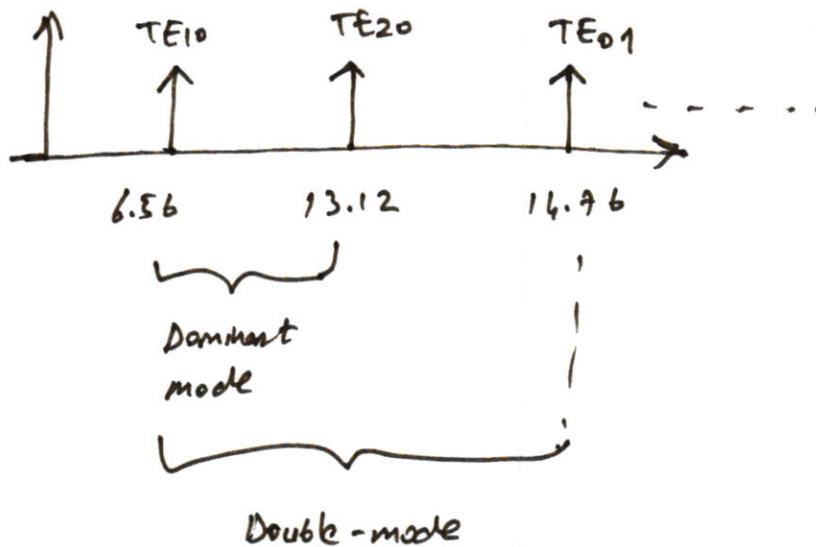
$$f_{c_{mn}} = \frac{c}{2\pi\sqrt{\epsilon_0\mu_0}} \sqrt{\left(\frac{m\pi}{a}\right)^2 + \left(\frac{n\pi}{b}\right)^2}$$

Hence,  $TE_{10}$  is the dominant mode for RWG (for  $a > b$ )

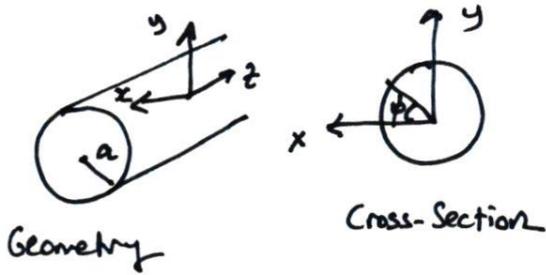
# For  $a = 2.286 \text{ cm}$ ,  $b = 1.016 \text{ cm}$

- i. Find first cut-off frequencies
- ii. Find frequency bandwidth of dominant mode
- iii. Find frequency bandwidth of double-mode region

Mode	$m$	$n$	$f_{c_{mn}}$ [GHz]
TE	1	0	6.562
TE	2	0	13.12
TE	0	1	14.76
TE <sub>1</sub> , TM	1	1	16.156
TE <sub>1</sub> , TM	1	2	30.248
TE <sub>1</sub> , TM	2	1	19.753



# CIRCULAR WAVEGUIDES



## TE MODES ( $E_z=0, H_z \neq 0$ )

$$(\nabla_t^2 + k_c^2) \psi = 0 \text{ where } \psi = H_z \text{ and } \frac{d\psi}{dn} = 0 \text{ on } \rho = a$$

$$\nabla_t^2 \psi = \nabla \cdot (\nabla \psi) = \nabla \cdot \left( \frac{d\psi}{d\rho} \vec{a}_\rho + \frac{1}{\rho} \frac{d\psi}{d\phi} \vec{a}_\phi + \frac{d\psi}{dz} \vec{a}_z \right)$$

$$\nabla_t^2 \psi = \frac{1}{\rho} \left[ \frac{d}{d\rho} \left( \rho \frac{d\psi}{d\rho} \right) + \frac{d}{d\phi} \left( \frac{1}{\rho} \frac{d\psi}{d\phi} \right) + \frac{d^2\psi}{dz^2} \right]$$

$$\psi = H_z(\rho, \phi, z) = h_z(\rho, \phi) e^{-i\beta z} \text{ and let's separate } h_z = R(\rho) P(\phi)$$

$$\nabla_t^2 h_z = \frac{d^2 h_z}{d\rho^2} + \frac{1}{\rho} \frac{dh_z}{d\rho} + \frac{1}{\rho^2} \frac{d^2 h_z}{d\phi^2}$$

$$= \rho \frac{d^2 R}{d\rho^2} + \rho \frac{dR}{d\rho} \cdot \frac{1}{\rho} + R \frac{1}{\rho^2} \frac{d^2 P}{d\phi^2} + k_c^2 R P = 0$$

$$\times \frac{1}{RP} \left\{ \begin{aligned} &= \frac{1}{R} \frac{d^2 R}{d\rho^2} + \frac{1}{R} \frac{1}{\rho} \frac{dR}{d\rho} + \frac{1}{P} \frac{1}{\rho^2} \frac{d^2 P}{d\phi^2} + k_c^2 = 0 \end{aligned} \right.$$

$$\times \rho^2 \left\{ \begin{aligned} &= \frac{\rho^2}{R} \frac{d^2 R}{d\rho^2} + \frac{\rho}{R} \frac{dR}{d\rho} + k_c^2 \rho^2 = -\frac{1}{P} \frac{d^2 P}{d\phi^2} \end{aligned} \right.$$

$$\text{OR } \left\{ \begin{aligned} &= \boxed{\rho^2 R'' + \rho R' + (\rho^2 k_c^2 - k_\phi^2) R = 0} \text{ where } k_\phi = \frac{-P''}{P} \end{aligned} \right.$$

↳ which is known as Bessel's Differential Equation

The solution is

$$R = C \cdot J_{k_\phi}(k_c \rho) + D \cdot N_{k_\phi}(k_c \rho)$$

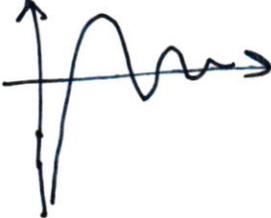
where  $J_\nu(x)$  is the Bessel function of the first kind of order  $\nu$  and

$N_\nu(x)$  is the Bessel function of the second kind of order  $\nu$ .

$$R \cdot P = H_z(\rho, \phi, z) = [C J_{k_\phi}(k_c \rho) + D N_{k_\phi}(k_c \rho)] [A \sin(k_\phi \phi) + B \cos(k_\phi \phi)] e^{-i\beta z}$$

since  $P'' + k_\phi^2 P = 0 \Rightarrow P = A \sin(k_\phi \phi) + B \cos(k_\phi \phi)$

★  $H_z(\rho, \phi, z) = H_z(\rho, \phi + 2\pi t, z)$  where  $t$  is an integer due to the circular symmetry. This can be true only if  $k_\phi$  becomes an integer. So,  $k_\phi = n$

★  $N_\nu(x)$  has the property  $\rightarrow$   as  $x \rightarrow 0, N_\nu(x) \rightarrow -\infty$

Since  $\rho=0$  is in the waveguide domain,  $N_\nu(x)$  cannot be a physical solution. Hence  $D \rightarrow 0$ . Hence;

$$h_z(\rho, \phi) = C \cdot [A \sin(n\phi) + B \cos(n\phi)] J_n(k_c \rho)$$

$\vec{n} = \vec{a}_y, \frac{dh_z}{dn} \Big|_{\rho=a} = 0 \Rightarrow \frac{dJ_n(k_c \rho)}{d\rho} = 0 \Rightarrow k_c = \frac{\rho'_{nm}}{a}$  where  $\rho'_{nm}$  is the  $n^{\text{th}}$  root of derivative  $n^{\text{th}}$  order Bessel of function

$$\beta_{nm} = \sqrt{k^2 - \left(\frac{\rho'_{nm}}{a}\right)^2}$$

$\rho'_{nm}$  values for some  $TE^{mn}$  modes in circular waveguide

$n \backslash m$	1	2	3
0	3.83	7.016	10.174
1	1.841	5.331	8.536
2	3.054	6.706	9.970

minimum value  
 $TE^{11}$  is the dominant mode in circular waveguide

## TM Modes

Similar derivations are valid. Hence

$$E_z(\rho, \phi, z) = [A \sin(n\phi) + B \cos(n\phi)] \cdot C \cdot J_n(k_c \rho) e^{-i\beta z}$$

$\varphi = 0$  on PEC where  $\varphi = E_z \rightarrow k_c = \frac{p_{nm}}{a}$  where  $p_{nm}$  is the  $m^{\text{th}}$  root of the  $n^{\text{th}}$  degree Bessel function  $J_n$

Some  $p_{nm}$  values

n \ m	1	2	3
0	2.4	5.52	8.65
1	3.83	7.01	10.17
2	5.13	8.417	11.62

$$\beta_{nm} = \sqrt{k^2 - \left(\frac{p_{nm}}{a}\right)^2}$$

Still, there exists no lower mode in TM modes than  $TE_{11}$  which is the dominant mode for circular waveguide.